

自然単位系

$$c = \hbar = 1$$

$$[\text{長さ}] = [\text{時間}]$$

$$[\text{質量}] = [\text{運動量}] = [\text{エネルギー}]$$

$$[\text{エネルギー}] = [\text{時間}]^{-1}$$

$$\psi(t) \sim e^{-iEt + ip \cdot x}$$

換算

$$c = 3 \times 10^{10} \text{ cm s}^{-1}$$

$$\hbar = 6.6 \times 10^{-25} \text{ GeV s}$$

$$\hbar c = 0.2 \text{ GeV fm}$$

$$(\hbar c)^2 = 0.39 \text{ GeV}^2 \text{ mb}$$

$$\begin{aligned} 1 \text{ fm} &= 10^{-15} \text{ m} \\ &= 10^{-13} \text{ cm} \end{aligned}$$

特殊相対論

4-vector $(a^0, a^1, a^2, a^3) = (a^0, \mathbf{a})$

内積 $a \cdot b = a^0 b^0 - \mathbf{a} \cdot \mathbf{b}$ Lorentz 不変

粒子の momentum

$(E, \mathbf{p}) = P^\mu$

$P^2 = E^2 - p^2 = m^2$: 質量

$\beta = \frac{|\mathbf{p}|}{E} = \frac{v}{c}$: 速度 $0 \leq \beta < 1$

$\gamma = \frac{1}{\sqrt{1-\beta^2}} = \frac{E}{m}$ Lorentz factor $1 \leq \gamma < \infty$

エネルギー $E = \gamma m$

運動量 $p = \beta E = \gamma \beta m$

粒子の崩壊

τ : lifetime $n(t) = n_0 e^{-t/\tau}$ (in rest frame)

走っている粒子は寿命が伸びる。

みかけの寿命 = $\gamma \tau$

走る距離 = $\gamma \beta c \tau$

(例) $p = 1 \text{ GeV}$ の K^+ , K_S^0 が興味の対象

$m_K \sim 0.5 \text{ GeV} \rightarrow \gamma \beta = \frac{p}{m} \sim 2$

$\tau_{K_S} = 0.9 \times 10^{-10} \text{ s} \rightarrow \gamma \beta c \tau \approx 5 \text{ cm}$

$\tau_{K^+} = 1.2 \times 10^{-8} \text{ s} \rightarrow \text{ " } 3 \text{ m}$

t/τ	$e^{-t/\tau}$
ϵ	$1 - \epsilon$
10^{-3}	0.9990
10^{-2}	0.990
10^{-1}	0.905
1	0.368
10	4×10^{-5}
100	4×10^{-44}

Why Collider ?

Beam - Fixed target

$$\begin{array}{ccc} \bullet \longrightarrow & \bullet & \\ P_1 = (E, 0, 0, E) & P_2 = (M, 0, 0, 0) & E \gg M \end{array}$$

$$S = (P_1 + P_2)^2 = (E+M)^2 - E^2 \approx 2ME$$

(重心系エネルギー²)

Beam - Beam

$$\begin{array}{ccc} \bullet \longrightarrow & \longleftarrow \bullet & \\ P_1 = (E, 0, 0, E) & P_2 = (E, 0, 0, -E) & \end{array}$$

$$S = (P_1 + P_2)^2 = 4E^2$$

重心系 = 実験室系

Asymmetric collider

(例: HERA)

e 30 GeV

p 820 GeV

$$\begin{array}{ccc} \circ \longrightarrow & \longleftarrow \bullet & \\ (E_1, 0, 0, E_1) & (E_2, 0, 0, -E_2) & \end{array}$$

$\sqrt{S} = 314 \text{ GeV}$

$$S = 4E_1E_2$$

Collider の弱点: Small rate

beam $\sim 10^{12}$ particles

solid target $\sim 10^{24}$ particles/cm³

Collider 実験の基本思想

あらゆる event をできるだけ記録する

→ さまざまな解析を同じデータに基づいて行う.

e^+e^- の event rate は低い

(例) $\sqrt{s} = 50 \text{ GeV}$ $\sigma(e^+e^- \rightarrow \text{hadron}) \sim 0.2 \text{ nb} = 2 \times 10^{-34} \text{ cm}^2$

$\mathcal{L} = 10^{31} \text{ cm}^2 \text{ s}^{-1} \rightarrow \text{event rate} = 2 \times 10^{-3} \text{ s}^{-1} \sim 10 \text{ 分} = 1 \text{ 回}$
(10^8 beam crossing = 1回)

ほとんどは $\underbrace{\text{ゴミ}}_{\text{の trigger}}$ (宇宙線 muon, beam-gas events)

Special purpose detectors も少数ながらある

ASP (PEP) : Single photon search

SHIP (TRISTAN) 他: Monopole search

加速器のエネルギー

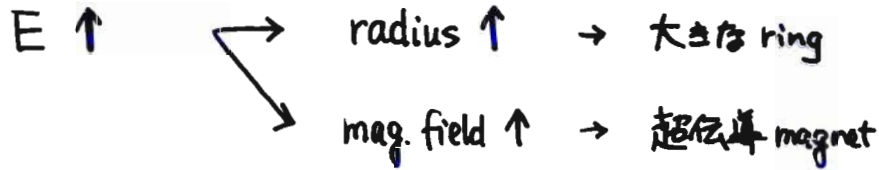
陽子 Synchrotron

最高エネルギーは最大磁場できまる

$$R = \frac{p}{eH}$$

$$\omega = \frac{eH}{E}$$

bending radius $\propto \frac{\text{momentum}}{\text{磁場}}$



電子 synchrotron

最高エネルギーは Synchrotron radiation できまる

$$I = \frac{2e^4 H^2 p^2}{3m^4}$$

energy loss $\propto \frac{E^4}{R^2}$



Storage ring のコスト $\propto R + R \cdot \frac{E^4}{R^2} \sim E^2$ (optimum) ($R \sim E^2$)

\uparrow 建設費 \uparrow To keep luminosity

\rightarrow Linear collider cost $\sim E$

Luminosity

実験の sensitivity のめやす

$$\mathcal{L} \sim \frac{N_{e^+} N_{e^-} f n_B}{A}$$

Luminosity ↑

N_{e^\pm} : Number of e^\pm per bunch

↑
増やす

f : Frequency (= $c / \text{一周の長さ}$)

n_B : バンチ数

↑

A : ビームの断面積

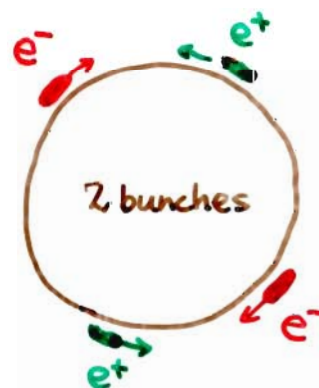
↓
小さくする

単位:

$$\text{cm}^{-2} \text{sec}^{-1}$$

Typical collider has

$$\sim 10^{30} \text{ cm}^{-2} \text{sec}^{-1}$$



Luminosity × cross section = 一秒あたりの反応数

Integrated Luminosity

$$\int \mathcal{L} dt$$

Luminosity の 時間積分

Integrated luminosity × cross section = 全反応数

単位

$$\text{pb}^{-1} = 10^{36} \text{ cm}^{-2}$$

$$\text{nb}^{-1} = 10^{33} \text{ cm}^{-2}$$

$$\mathcal{L} \sim 10^{31} \text{ cm}^{-2} \text{ s}^{-1} \times 2$$

年間 ($\sim 10^7 \text{ s}$)

$$\rightarrow 10^{38} \text{ cm}^{-2} = 100 \text{ pb}^{-1}$$

Cross section

断面積 : 反応の起こりやすさ

次元: $[長さ]^2$
or $[エネルギー]^{-2}$

単位 cm^2 $1b = 10^{-24} cm^2$

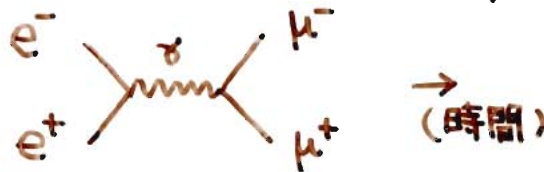
1 mb	$10^{-27} cm^2$	← hh
1 μb	10^{-30}	← γh
1 nb	10^{-33}	
1 pb	10^{-36}	← TRISTAN
1 fb	10^{-39}	← TeV e^+e^-

$$\sigma_{pp, tot} \sim 40 \text{ mb} \quad (\text{at minimum})$$

$$\sigma(e^+e^- \rightarrow \mu^+\mu^-) = \frac{4\pi\alpha^2}{3s} = \frac{87 \text{ nb}}{s(\text{GeV}^2)} \quad \text{: unit of } R$$

(QED lowest order)

$\sqrt{s} = \text{c.m. energy}$



e^+e^- Colliding Facilities

Name	Location	Typical Energy	Major Detectors	Started
SPEAR	SLAC (Stanford)	$J/\psi - 7.4$	{ Mark I Mark II DELCO Crystal Ball ^o Mark III	1973
DORIS	DESY (Hamburg)	$J/\psi - 7.$	{ PLUTO DASP LENA ^o	1974
DORIS II		Υ region	ARGUS Crystal Ball	↓ 1978
VEPP-4	ITEP (Novosibirsk)	Υ region		
CESR	Cornell (Ithaca)	Υ region	CLEO CUSP ^o	1980
PEP	SLAC	29	{ Mark II MAC HRS DELCO TPC-2y ASP	1980
PETRA	DESY	34 - 47	{ PLUTO Mark J TASSO JADE CELLO	1978 - 86
TRISTAN	KEK (Tsukuba)	50 - 57 ↑	VENUS TOPAZ AMY	1986
SLC	SLAC	Z	Mark II SLD	1988 ?
LEP	CERN (Genève)	Z Phase II: 190	ALEPH DELPHI L3 OPAL	1989
CLIC		TeV		
JLC				
BEPC	北京	ψ region		
"B factory"	Several proposals	Υ region		

Current and Future Colliders

e^+e^-	location	c.m. energy	luminosity	start	detectors
DAΦNE	Frascati (Italy)	ϕ region	$(1 \rightarrow 5) 10^{32}$	1997	
BEPC	北京 (中国)	ψ region	10^{31}	1989	BES
CESR	Cornell (USA)	Υ region	$5 \cdot 10^{32} \uparrow$	1979	CLEO
KEK B	つくば	"	10^{34}	1999	BELLE
PEP II	SLAC (USA)	"	$3 \cdot 10^{33}$	1999	BABAR
SLC	SLAC	Z region	10^{30}	1989	SLD
LEP	CERN	"	$2 \cdot 10^{31}$	1989	} ALEPH DELPHI L3 OPAL
LEP II		171 \rightarrow 182 GeV \rightarrow		1996	
\times LC	?	500 \rightarrow TeV region	$10^{33} - 10^{34}$	200?	

$p\bar{p}$

Tevatron	Fermilab (USA)	1.8 TeV \uparrow	$2 \cdot 10^{31} \uparrow$	1987	CDF DØ
LHC	CERN	14 TeV	10^{34}	2004	ATLAS CMS

e^+p

HERA	DESY (Germany)	300 GeV ($e^+ 27$ $p 820$)	10^{31}	1992	H1 ZEUS
------	----------------	------------------------------------	-----------	------	------------

Particle Lifetimes

γ, e, p, ν 安定

n 10^3 sec

μ 2.2×10^{-6} sec

lifetime ∞

K_L 5.2×10^{-8}

π^\pm 2.6×10^{-8}

K^\pm 1.2×10^{-8}

Λ 2.6×10^{-10}

K_S 0.9×10^{-10}

$p\pi^-$ 64%

$\pi^+\pi^-$ 69%

$\Sigma^+, \Sigma^-, \Xi^0, \Xi^-, \Omega^-$ $(0.8 - 3) \times 10^{-10}$

B 1×10^{-12}

charm $(1 - 10) \times 10^{-13}$

$D^{*+} \rightarrow D^0\pi^+$
 $\rightarrow K^-\pi^+$

τ 3×10^{-13}

π^0 0.9×10^{-16}

$\gamma\gamma$ 99%

lifetime 0

η 6×10^{-19}

Σ^0 6×10^{-20}

J/ψ 1×10^{-20}

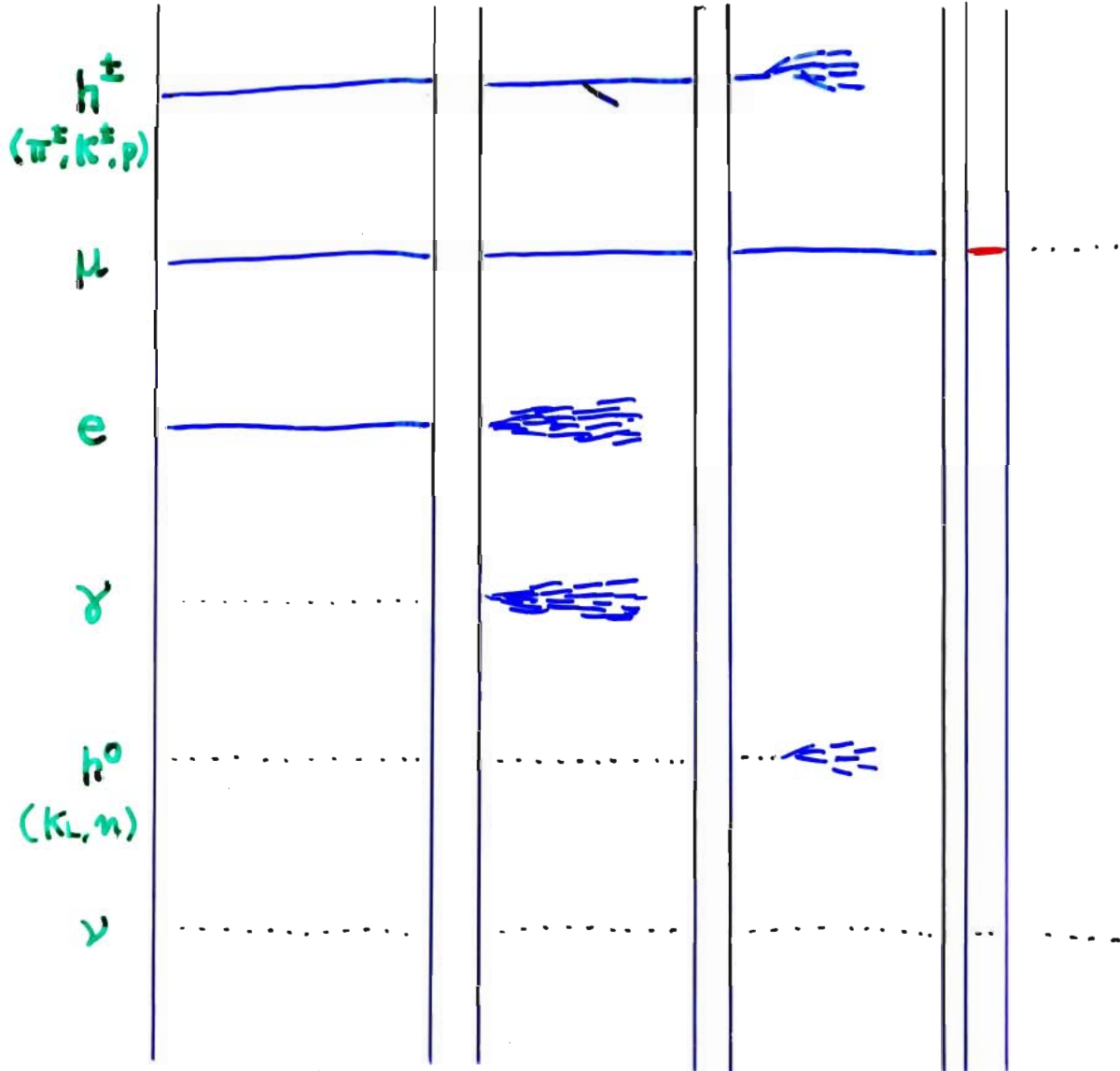
$e^+e^-, \mu^+\mu^-$ 7% each

\vdots

ρ 4×10^{-24}

W, Z $(2 - 3) \times 10^{-25}$

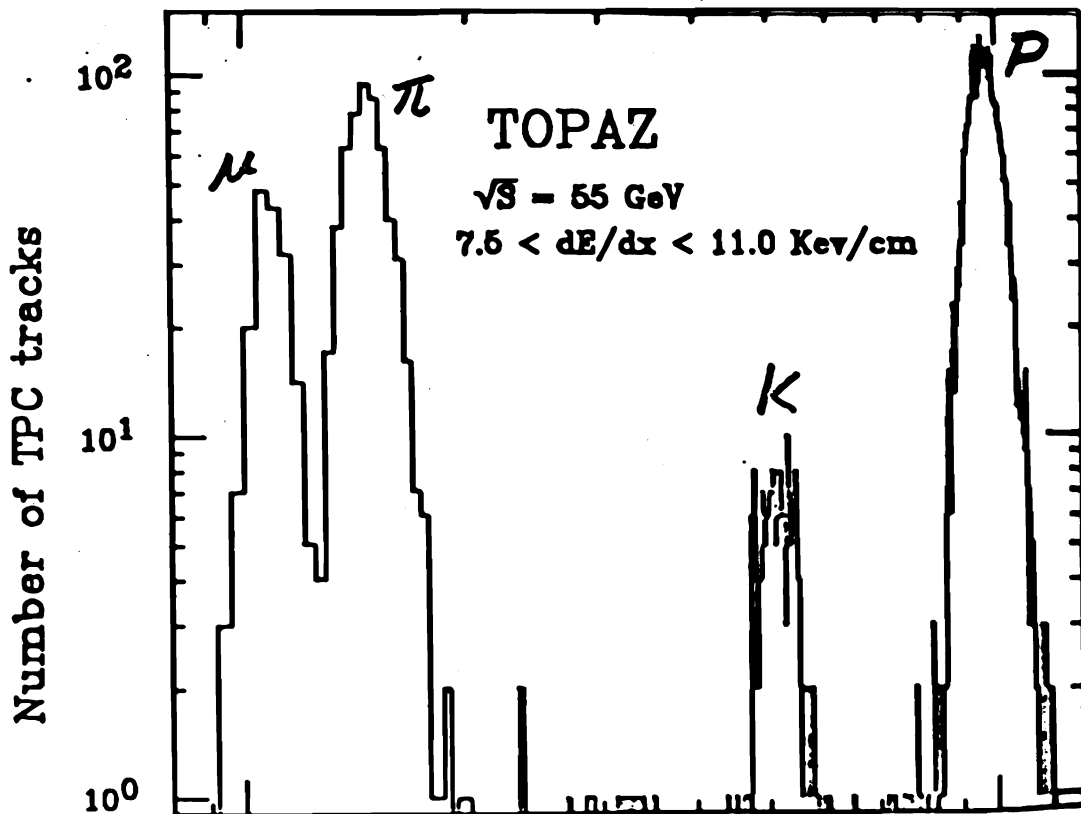
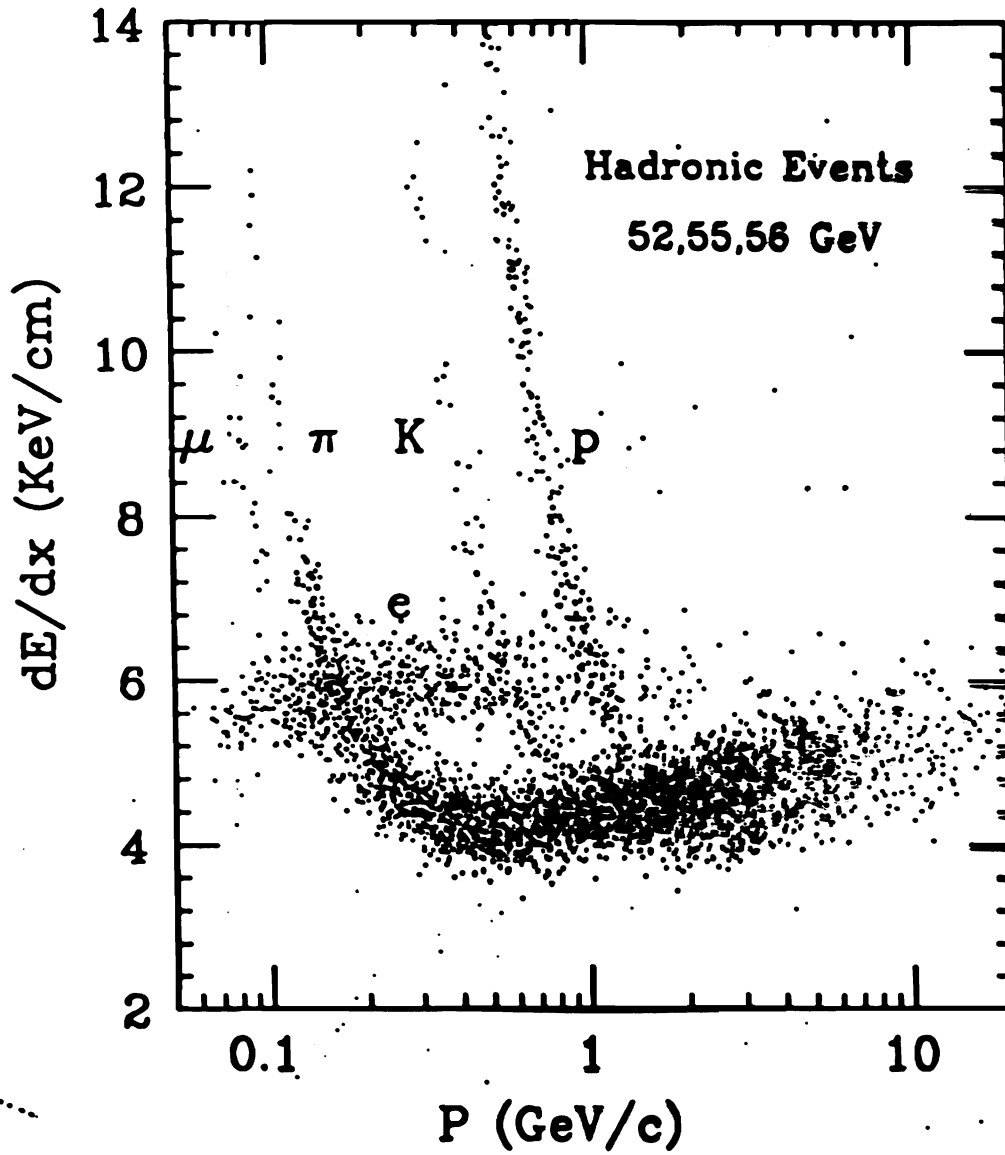
Particle identification



$\pi/K/p$ separation : measure p and β

- TOF (time of flight)
- $\frac{dE}{dx}$
- Čerenkov

Particle ID by TOPAZ-TPC



Radiative corrections : Where important ?

Electroweak radiative corrections : $O(\alpha)$ in general

But, large corrections appear with a log factor

$$\propto \log(\text{大数})$$

in the following cases

(1) Related to **ultraviolet** divergence

$$\propto \log \frac{\Lambda^2}{q^2} \xrightarrow[\langle 1 \rangle = \cancel{+}]{\text{on-shell}} \propto \left(\log \frac{\Lambda^2}{q^2} - \log \frac{\Lambda^2}{m^2} \right) = -\propto \log \frac{q^2}{m^2}$$



\Rightarrow use running coupling

(2) Related to **soft (infrared)** divergence $\leftrightarrow m_\gamma = 0$

$$\propto \log \frac{S}{E_{\min}^2}$$

divergence cancel between real and virtual corrections



No large correction for "total" cross section

(3) Related to **collinear** singularity $\leftrightarrow m_e \rightarrow 0$

(3a) **initial radiation**



$$\propto \log \frac{S}{m_e^2}$$

effective reduction of "c.m. energy"
absorbable into "structure function."

(3b) **final radiation**



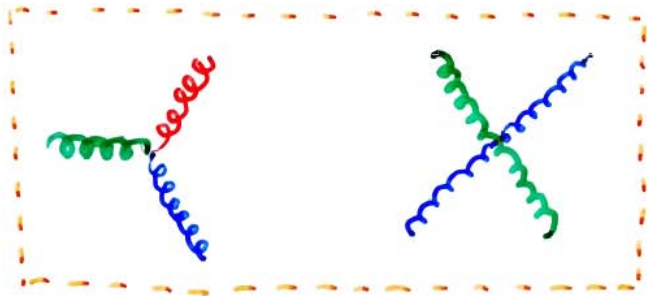
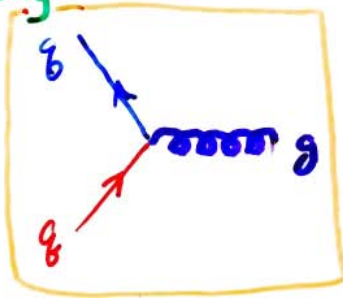
$$\propto \log \theta_{\min}$$

singularity cancel between real and virtual

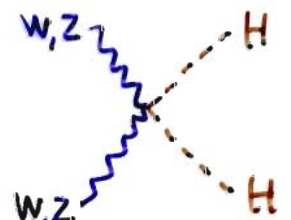
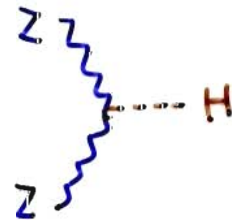
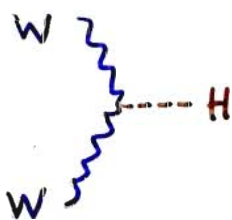
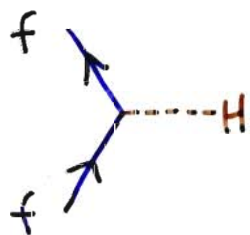
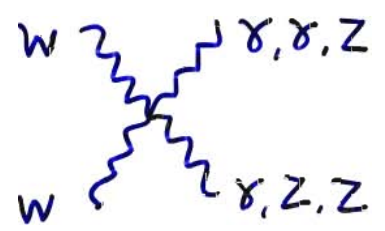
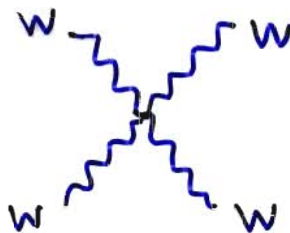
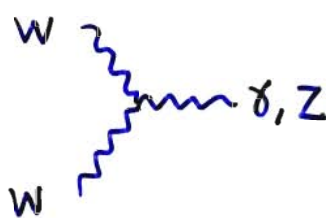
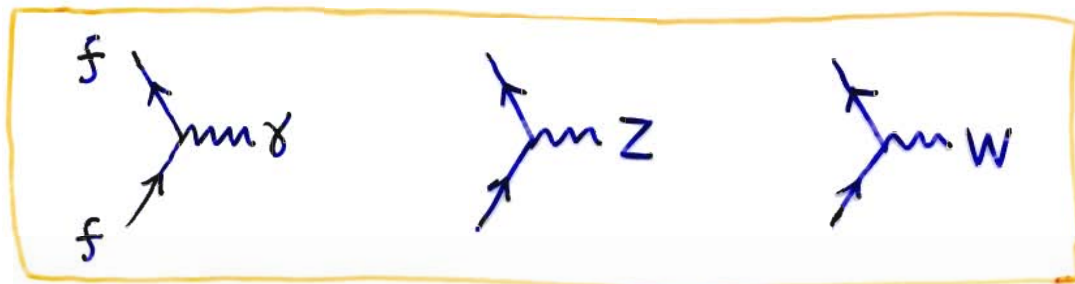
no large correction for "total" cross section

Standard model interactions

Strong



Electroweak



Standard model parameters

1. Basic parameters

Parameters in the fundamental Lagrangian

Should be measured by experiments in principle

		number
Gauge couplings	g_s, g, g'	3
Yukawa couplings	$f^{(d)}, f^{(l)}, h^{(u)}$	13 (3 generations)
Higgs self coupling	λ	1
Higgs mass Term	μ^2	1

Alternatively one may take

$\alpha, G_F, (\sin^2\theta_W \text{ or } M_Z \text{ or } M_W)$

$m_e, m_g, \text{ Kobayashi-Maskawa matrix}$

m_H

$\Lambda_{QCD} \text{ (or } \alpha_s)$

2. Secondary parameters

Calculable in principle, not in practice

f_π, f_K, \dots (decay constant)

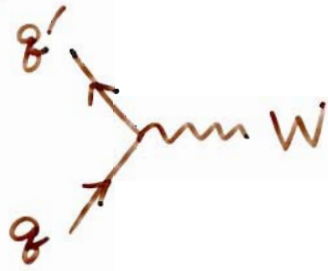
hadron masses, magnetic moments etc.

hadron structure functions

various matrix elements

Kobayashi-Maskawa matrix

Charged current interaction of quarks



$$\mathcal{L} = \frac{-g}{2\sqrt{2}} (\bar{u} \ \bar{c} \ \bar{t}) \gamma_{\mu} (1 - \gamma_5) V \begin{pmatrix} d \\ s \\ b \end{pmatrix} \cdot W^{+\mu} + \text{h.c.}$$

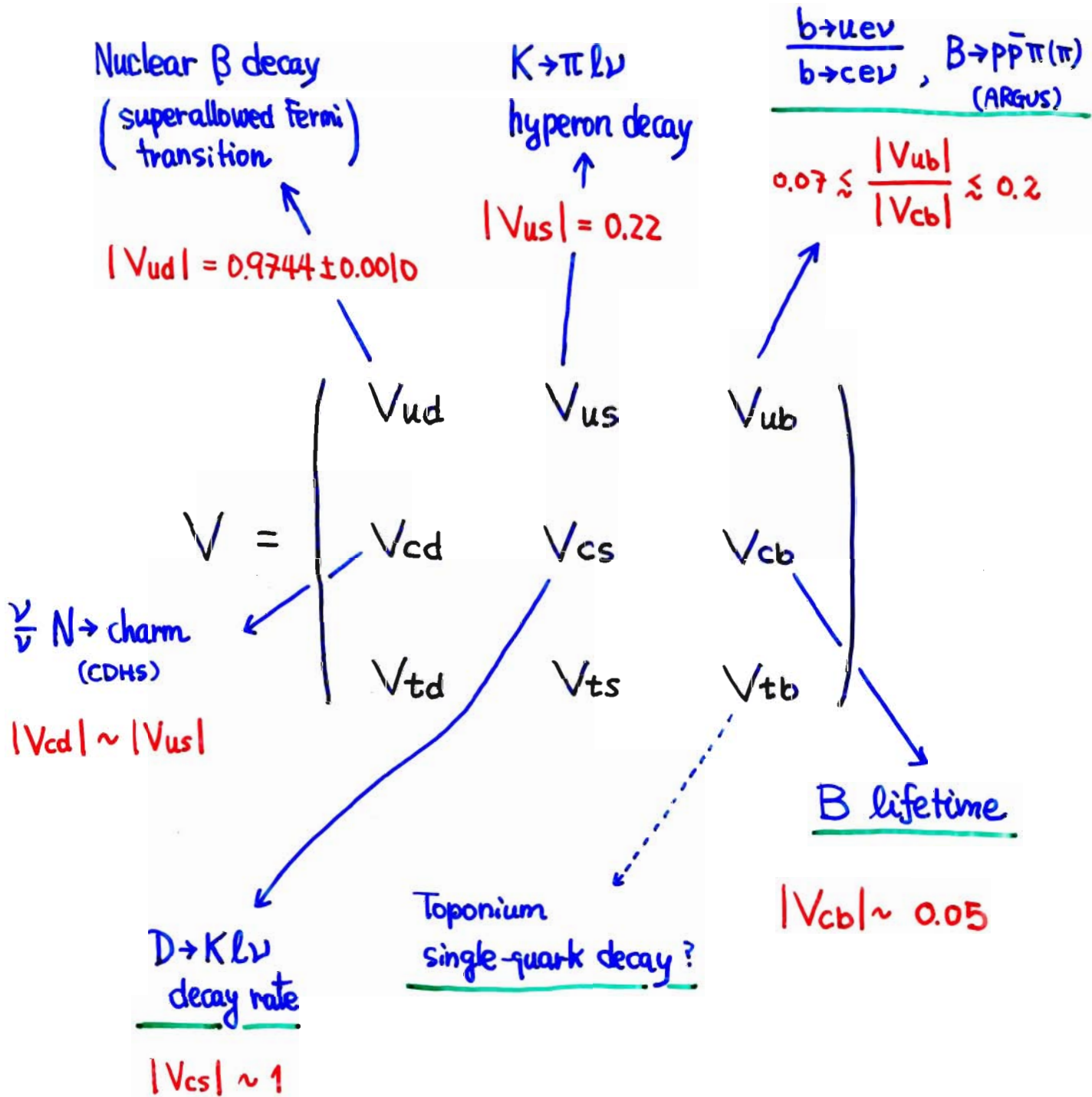
\swarrow
3x3 unitary matrix

V : 9成分 \rightarrow quark fields の relative phase
は無意味な γ (6-1) = 5成分は
unphysical

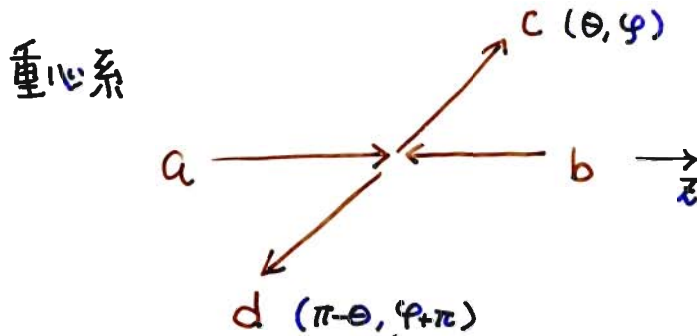
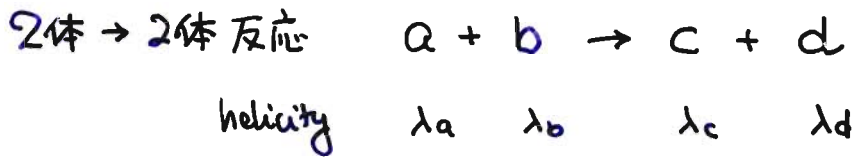
\rightarrow 4 physical degrees of freedom

$\left. \begin{array}{l} 3 \text{ angles} \\ 1 \text{ phase (CP-violation)} \end{array} \right\}$

Measuring KM Matrix



Helicity amplitude



Jacob-Wick formalism : 部分波展開

$$S_{fi} = \delta_{fi} + i(2\pi)^4 \delta^4(P_f - P_i) M_{fi}$$

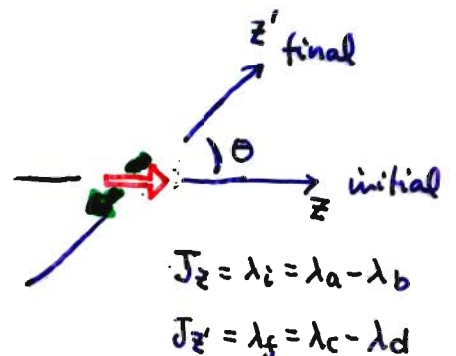
$$M_{fi} = \frac{8\pi}{\sqrt{\bar{\beta}_i \bar{\beta}_f}} \sum_{J=0}^{\infty} (2J+1) T_{\lambda_a \lambda_b; \lambda_c \lambda_d}^J(E_{cm}) d_{\lambda_i \lambda_f}^J(\theta) e^{i(\lambda_i - \lambda_f)\varphi}$$

$$\lambda_i = \lambda_a - \lambda_b$$

$$\lambda_f = \lambda_c - \lambda_d$$

$$\bar{\beta}_i = \frac{2|P_i|}{\sqrt{s}} \quad \bar{\beta}_f = \frac{2|P_f|}{\sqrt{s}}$$

$d_{\lambda_i \lambda_f}^J(\theta)$: Wigner's d function
(回転行列)



helicity amplitude は

$$M_{fi} = \tilde{M}_{\lambda_a \lambda_b; \lambda_c \lambda_d}(\sqrt{s}, \cos\theta) \underbrace{d_{\lambda_i \lambda_f}^{j_0}(\theta)}_{\text{「最低限の」角分布}} e^{i(\lambda_i - \lambda_f)\varphi}$$

と書ける。 $j_0 = \max(|\lambda_i|, |\lambda_f|)$

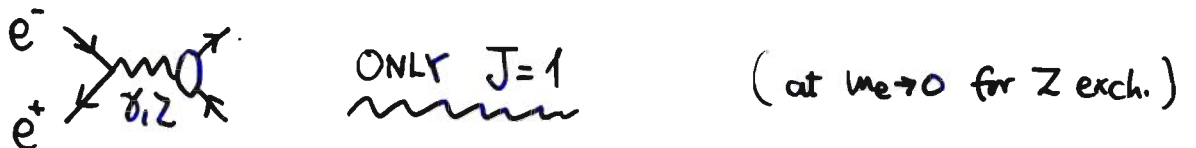
「最低限の」角分布

$$\left(\begin{array}{l} \text{但し} \\ \tilde{M} = \frac{8\pi}{\sqrt{\beta_i \beta_f}} \sum_{J=j_0}^{\infty} (2J+1) T_{\lambda_a \lambda_b; \lambda_c \lambda_d}^J(\sqrt{s}) \cdot \frac{d_{\lambda_i \lambda_f}^J(\theta)}{d_{\lambda_i \lambda_f}^{j_0}(\theta)} \end{array} \right)$$

\downarrow
 $\cos\theta$ a polynomial

S-channel exchange の場合 J は 特定の値しかとらない

(例) $e^+ e^-$. one photon exchange (Z)



$$M_{fi} = \tilde{M}(\sqrt{s}) d_{\lambda_i \lambda_f}^1(\theta) e^{i(\lambda_i - \lambda_f)\varphi}$$

d 関数

$$d_{mm'}^j(\theta) = [(j+m)!(j-m)!(j+m')!(j-m')!]^{1/2} \\ \times \sum_k \frac{(-1)^k}{(j-m'-k)!(j+m-k)!(k+m'-m)!k!} \\ \times \left(\cos\frac{\theta}{2}\right)^{2j+m-m'-2k} \left(-\sin\frac{\theta}{2}\right)^{m'-m+2k}$$

$$d_{m'm}^j(\theta) = d_{mm'}^j(-\theta)$$

$$d_{mm'}^j(\pi-\theta) = (-1)^{j+m} d_{m,-m'}^j(\theta)$$

$$d_{mm'}^j(0) = \delta_{mm'}$$

$$d_{mm'}^j = (-1)^{m-m'} d_{m'm}^j = (-1)^{m-m'} d_{-m,-m'}^j = d_{-m',-m}^j$$

$$\int_0^\pi d\cos\theta d_{mm'}^j(\theta) d_{mm'}^{j'}(\theta) = \frac{2}{2j+1} \delta_{jj'}$$

$j_0 \equiv \max(|m|, |m'|)$ とおくと

$$d_{mm'}^j(\theta) = d_{m_0 m_0}^{j_0}(\theta) \times (\cos\theta \text{ の多項式})$$

$$d_{\lambda\mu}^{1/2}(\theta)$$

λ	μ	$+1/2$	$-1/2$
$+1/2$		$\cos \frac{\theta}{2}$	$-\sin \frac{\theta}{2}$
$-1/2$		$\sin \frac{\theta}{2}$	$\cos \frac{\theta}{2}$

$$d_{\lambda\mu}^1$$

λ	μ	$+1$	0	-1
$+1$		$\frac{1+\cos\theta}{2}$	$-\frac{1}{\sqrt{2}}\sin\theta$	$\frac{1-\cos\theta}{2}$
0		$\frac{1}{\sqrt{2}}\sin\theta$	$\cos\theta$	$-\frac{1}{\sqrt{2}}\sin\theta$
-1		$\frac{1-\cos\theta}{2}$	$\frac{1}{\sqrt{2}}\sin\theta$	$\frac{1+\cos\theta}{2}$

$$d_{00}^0 = 1$$

$d_{\lambda\mu}^2$

λ	μ	2	1	0	-1	-2
2		$\left(\frac{1+\cos\theta}{2}\right)^2$	$-\frac{1+\cos\theta}{2}\sin\theta$	$\frac{\sqrt{6}}{4}\sin^2\theta$	$-\frac{1-\cos\theta}{2}\sin\theta$	$\left(\frac{1-\cos\theta}{2}\right)^2$
1		$\frac{1+\cos\theta}{2}\sin\theta$	$\frac{1+\cos\theta}{2}(2\cos\theta-1)$	$-\sqrt{\frac{3}{2}}\sin\theta\cos\theta$	$\frac{1-\cos\theta}{2}(2\cos\theta+1)$	$-\frac{1-\cos\theta}{2}\sin\theta$
0		$\frac{\sqrt{6}}{4}\sin^2\theta$	$\sqrt{\frac{3}{2}}\sin\theta\cos\theta$	$\frac{3}{2}\cos^2\theta - \frac{1}{2}$	$-\sqrt{\frac{3}{2}}\sin\theta\cos\theta$	$\frac{\sqrt{6}}{4}\sin^2\theta$
-1		$\frac{1-\cos\theta}{2}\sin\theta$	$\frac{1-\cos\theta}{2}(2\cos\theta+1)$	$\sqrt{\frac{3}{2}}\sin\theta\cos\theta$	$\frac{1+\cos\theta}{2}(2\cos\theta-1)$	$-\frac{1+\cos\theta}{2}\sin\theta$
-2		$\left(\frac{1-\cos\theta}{2}\right)^2$	$\frac{1-\cos\theta}{2}\sin\theta$	$\frac{\sqrt{6}}{4}\sin^2\theta$	$\frac{1+\cos\theta}{2}\sin\theta$	$\left(\frac{1+\cos\theta}{2}\right)^2$

Helicity amplitudes ($\varphi=0$)

$e^-e^+ \rightarrow \mu^-\mu^+$: nonzero amplitudes for $m_e, m_\mu \rightarrow 0$

$$s-\gamma \quad \mathcal{M} = -2e^2 d^1 \quad |\lambda_i| = |\lambda_f| = 1$$

$$s-Z \quad -2g_z^2 (V_e - \lambda_i a_e)(V_\mu - \lambda_f a_\mu) \frac{s}{s-m_Z^2} d^1$$

$$|\lambda_i| = |\lambda_f| = 1$$

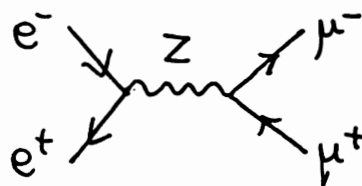
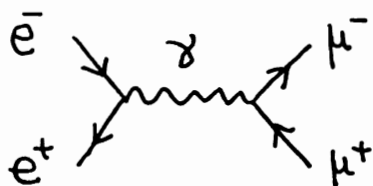
$$g_z = e / \sin\theta_w \cos\theta_w$$

$$V_e = V_\mu = -\frac{1}{4} + \sin^2\theta_w \ll a_e, a_\mu$$

$$a_e = a_\mu = -\frac{1}{4}$$

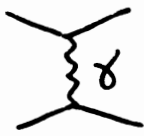
$$d_{11}^1 = d_{-1,-1}^1 = \frac{1}{2}(1 + \cos\theta)$$


$$d_{1,-1}^1 = d_{-1,1}^1 = \frac{1}{2}(1 - \cos\theta)$$



$e^-e^+ \rightarrow e^-e^+$ nonzero amplitudes for $m_e \rightarrow 0$

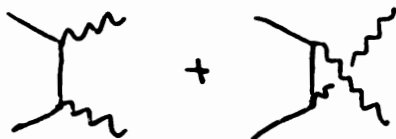
s- γ , s- Z : same as for $e^-e^+ \rightarrow \mu^-\mu^+$ ^{with} ($V_\mu, a_\mu \rightarrow V_e, a_e$)

t- γ	$M = \frac{4e^2}{1-\cos\theta} d^1$	$\lambda_i = \lambda_f = \pm 1$	$\begin{matrix} h\bar{h} & h'\bar{h}' \\ (+-) & (+-) \\ (-+) & (-+) \end{matrix}$
	$= \frac{4e^2}{1-\cos\theta}$	$\lambda_i = \lambda_f = 0$ and $h = h'$	$\begin{matrix} (++) & (++) \\ (--) & (--) \end{matrix}$

t- Z	$M = \frac{4g_z^2}{1 + \frac{2m_z^2}{s} - \cos\theta} (V_e - \lambda_i a_e)^2 d^1$	$\lambda_i = \lambda_f = \pm 1$	
	$= \frac{4g_z^2}{1 + \frac{2m_z^2}{s} - \cos\theta} (V_e^2 - a_e^2)$	$\lambda_i = \lambda_f = 0$ and $h = h'$	

$e^-e^+ \rightarrow \gamma\gamma$

t-e \oplus u-e	$M = -\frac{4e^2}{1-\cos^2\theta} d^2$	$ \lambda_i = 1, \lambda_f = 2$
------------------	--	------------------------------------



Unitarity and cross section

$$S^\dagger S = 1 \quad (\text{確率の保存})$$

for partial wave amplitudes

$$\text{elastic channel } (f=i) \quad |T^J| \leq 2$$

$$\text{inelastic channel} \quad |T^J| \leq 1$$

Integrated cross section

$$\sigma = \frac{4\pi}{S \beta_i^2} \sum_{J=0}^{\infty} (2J+1) |T^J|^2$$

$e^+e^- \xrightarrow{\gamma^*}$ final の場合

$J=1$ only

$$T^{J=1} \sim e^2$$

$$\Rightarrow \sigma \sim \frac{\alpha^2}{S}$$

Unit of R

$$\sigma(e^+e^- \rightarrow \mu^+\mu^-) = \frac{4\pi\alpha^2}{3S}$$